

A SPECTRAL EXPANSION FOR SCHRÖDINGER OPERATOR

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Abstract

In this paper we consider the Schrödinger operator L generated in $L^2(\mathbf{R}_+)$ by

$$y'' + q(x)y = \mu y, \quad x \in \mathbf{R}_+ := [0, \infty)$$

subject to the boundary condition

$$y'(0) - hy(0) = 0,$$

where, q is a complex valued function summable in $[0, \infty$ and $h \neq 0$ is a complex constant, μ is a complex parameter. We have assumed that

$$\sup_{x \in \mathbf{R}_+} \{\exp(\varepsilon\sqrt{x}) |q(x)|\} < \infty, \quad \varepsilon > 0,$$

holds which is the minimal condition that the eigenvalues and the spectral singularities of the operator L are finite with finite multiplicities. Under this condition we have given the spectral expansion formula for the operator L using an integral representation for the Weyl function of L . Moreover we also have investigated the convergence of the spectral expansion.

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Key Words : *Spectrum, Weyl Function, Spectral Expansion.*

1. Introduction

Let L denote the operator generated in $L^2(\mathbf{R}_+)$ by differential equation

$$(1.1) \quad -y'' + q(x)y = \mu y, \quad x \in \mathbf{R}_+$$

and the boundary condition

$$(1.2) \quad y'(0) - hy(0) = 0,$$

here q is a complex valued function, and $h \neq 0$ is a complex number, μ is a complex parameter. The spectral analysis of L was first investigated by Naimark [6]. In this paper, he has proved that some of the poles of the resolvent's kernel of L are not the eigenvalues of the operator. Moreover he has proved that these poles are on the continuous spectrum. (Schwartz named these poles as spectral singularities [7]. Furthermore Naimark has shown that if the condition

$$(1.3) \quad \int_0^{\infty} e^{\varepsilon x} |q(x)| dx < \infty, \quad \varepsilon > 0$$

holds, then L has a finite number of eigenvalues and spectral singularities each of them is of finite multiplicity.

The effect of spectral singularities in the spectral expansion of the operator L in terms of principal functions under the condition (1.3) has been investigated in [3].

The spectral expansion in terms of principal vectors for the non-selfadjoint discrete Dirac operator has been given in [1] using the integral representation for the Weyl function of the operator.

Suppose that the condition

$$(1.4) \quad \sup_{x \in \mathbf{R}_+} \{\exp(\varepsilon\sqrt{x}) |q(x)|\} < \infty, \quad \varepsilon > 0,$$

holds. Under the condition (1.4) the operator L has a finite number of eigenvalues and spectral singularities each of them with finite multiplicity (see [2]).

In this paper, using the similar technique used in [1], we have obtained the spectral expansion of the operator L using an integral representation for the Weyl function of L under the condition (1.4). Moreover we have investigated the convergence of the spectral expansion.

2. Preliminaries

We denote the solutions of (1.1) satisfying the initial conditions

$$(2.1) \quad \begin{aligned} w(0, \sqrt{\mu}) &= 1, \quad w'(0, \sqrt{\mu}) = h, \\ \varphi(0, \mu) &= 1, \quad \varphi'(0, \mu) = 0, \end{aligned}$$

$$\theta(0, \mu) = 0, \quad \theta'(0, \mu) = -1$$

by $w(x, \sqrt{\mu})$, $\varphi(x, \mu)$, $\theta(x, \mu)$, respectively, the set of all finite functions in $L^2(\mathbf{R}_+)$ by C_0^∞ and the w -Fourier transform of an arbitrary element $f \in C_0^\infty$ by $E_f(\sqrt{\mu}) = \int_0^\infty f(x) w(x, \sqrt{\mu}) dx$. Moreover \mathbf{C} denotes the complex plane. Marchenko has obtained an analytical expression of the generalized spectral function of the operator L by using Weyl function [4].

We give some important results which we need.

Theorem 2.1 ([4]). *Suppose the potential of the equation (1.1) satisfies the condition*

$$(2.2) \quad r_0 \leq \operatorname{Im} q \leq r_1.$$

Then (1.1) has a unique solution given by

$$(2.3) \quad \psi(x, \mu) = \theta(x, \mu) + m(\mu) w(x, \sqrt{\mu}) \in L^2(\mathbf{R}_+)$$

for all μ which is not in the strip

$$(2.4) \quad r_0 \leq \operatorname{Im} \mu \leq r_1,$$

where $m(\mu)$ is analytic outside the strip (2.4).

The function $m(\mu)$ given in Theorem 2.1 is the Weyl function, and the function $\psi(x, \mu)$ is the Weyl solution of the operator L . Weyl function and Weyl solution may be obtained as follows:

$$(2.5) \quad \psi_1(x, \mu) = \frac{\psi(x, \mu)}{1 + hm(\mu)}, \quad m_1(\mu) = \frac{m(\mu)}{1 + hm(\mu)},$$

moreover $m(\mu)$ satisfies

$$(2.6) \quad \lim_{|\mu| \rightarrow \infty} m(\mu) = 0$$

outside the strip (2.4) ([4]).

Theorem 2.2 ([4]). *If the condition (2.2) holds, then for all $f \in C_0^\infty$, the generalized spectral function R of the operator L is given by the formula*

$$(2.7) \quad (R, E_f(\sqrt{\mu})) = \frac{1}{2\pi i} \lim_{\delta \rightarrow 0} \left\{ \int_{-\infty+i(r_1+\varepsilon)}^{\infty+i(r_1+\varepsilon)} e^{-\delta\mu^2} m_1(\mu) E_f(\sqrt{\mu}) d\mu \right. \\ \left. + \int_{\infty+i(r_0-\varepsilon)}^{-\infty+i(r_0-\varepsilon)} e^{-\delta\mu^2} m_1(\mu) E_f(\sqrt{\mu}) d\mu \right\} \\ - \sum [m_1(\mu) E_f(\sqrt{\mu})],$$

where $\varepsilon > 0$ is a sufficiently small number and residues are computed for all poles of $m_1(\mu)$ which is not in the strip $r_0 - \varepsilon \leq \text{Im } \mu \leq r_1 + \varepsilon$.

It is known that to every boundary value problem (1.1) – (1.2) there corresponds a functional $R \in Z'$, $f, g \in C_0^\infty$, such that

$$(2.8) \quad (R, E_f(\sqrt{\mu}) E_g(\sqrt{\mu})) = \int_0^\infty f(x) g(x) dx,$$

where Z' denotes the space of all linear, continuous functionals defined on Z which is the set of all even, entire functions of exponential type, summable on \mathbf{R} . (2.8) is called Marchenko-Parseval equality, and R is the generalized spectral function of L .

Theorem 2.3. ([5]). *Suppose that the condition*

$$(2.9) \quad \int_0^\infty x |q(x)| dx < \infty$$

holds. If $\mu = \lambda^2$ then the equation (1.1) has the solution

$$(2.10) \quad e(x, \lambda) = e^{i\lambda x} + \int_x^\infty K(x, t) e^{i\lambda t} dt$$

which is analytic in $\text{Im } \lambda > 0$ and continuous on $\text{Im } \lambda \geq 0$, satisfying $\lim_{x \rightarrow \infty} e(x, \lambda) e^{-i\lambda x} = 1$ for $\text{Im } \lambda \geq 0$. Furthermore $K(x, t)$ satisfies the inequality

$$|K(x, t)| \leq \frac{1}{2} \sigma\left(\frac{x+t}{2}\right) \exp \sigma_1(x),$$

where $\sigma(x) = \int_x^\infty |q(t)| dt$, $\sigma_1(x) = \int_x^\infty t |q(t)| dt$.

Now let us denote the eigenvalues and the spectral singularities and continuous spectrum of L by $\sigma_d(L)$, $\sigma_{ss}(L)$ and $\sigma_c(L)$, respectively.

$$\text{Clearly } \sigma_d(L) = \{\mu : \mu = \lambda^2, \text{Im } \lambda > 0, d(\lambda) = 0\},$$

$$\sigma_{ss}(L_\lambda) = \{\mu : \mu = \lambda^2, \lambda \in \mathbf{R}, d(\lambda) = 0\} \setminus \{0\},$$

$$\sigma_c(L) = [0, \infty,$$

where

$$(2.11) \quad d(\lambda) = e_x(0, \lambda) - h e(0, \lambda),$$

[6]. Therefore we have the following.

Theorem 2.4. *If the condition (1.4) holds then the operator L has a finite number of eigenvalues and spectral singularities with finite multiplicities.*

For the proof see for ex. [2].

Let us introduce the following notations

$$(2.12) \quad \widehat{e}(x, \mu) = e(x, \sqrt{\mu}), \quad \widehat{d}(\mu) = d(\sqrt{\mu}),$$

$\mu \in \mathbf{C} \setminus \mathbf{R}_+$ (i.e.: the complex plane except for nonnegative real numbers).

Remark 2.5. *If the condition (1.4) holds, then $|q(x)|$ will be bounded. From Theorem 2.1 Weyl's limit point case is satisfied for the operator L . So all solutions of (1.1) belonging to $L^2(\mathbf{R}_+)$ for μ outside the strip (2.4) are linearly dependent with the solution $\psi(x, \mu)$. From (2.5) $\psi_1(x, \mu) \in L^2(\mathbf{R}_+)$ for μ outside the strip (2.4) and $m(\mu) \neq \frac{-1}{h}$. Therefore considering (2.10) and (2.12) we arrive at $\widehat{e}(x, \mu) \in L^2(\mathbf{R}_+)$ for $\mu \in \mathbf{C} \setminus \mathbf{R}_+$. So there exists a constant $c \neq 0$ such that*

$$(2.13) \quad \psi_1(x, \mu) = c \widehat{e}(x, \mu),$$

moreover we get

$$(2.14) \quad m_1(\mu) = \frac{-\widehat{e}(\mu)}{\widehat{d}(\mu)}$$

by (2.3) and (2.13), where $\widehat{e}(\mu) = \widehat{e}(0, \mu)$, for μ outside the strip (2.4) and $m(\mu) \neq \frac{-1}{h}$.

Remark 2.6. If the condition (1.4) holds, then we can omit the summing factor $e^{-\delta\mu^2}$ in (2.7). In this case the function $m_1(\mu)$ is analytic and bounded in $\mathbf{C} \setminus \mathbf{R}_+$ except for the finite number of poles. So the formula (2.7) may be reduced into the form ([1])

$$(2.15) \quad (R, E_f(\sqrt{\mu})) = -\frac{1}{2\pi i} \int_{\gamma} m_1(\mu) E_f(\sqrt{\mu}) d\mu,$$

where γ is the contour enclosing the spectrum of the operator L , and $E_f(\sqrt{\mu})$ is the w -Fourier transform of $f \in C_0^\infty$. Taking $E_g(\sqrt{\mu})$ as the w -Fourier transform of $g \in C_0^\infty$. We know that Parseval's equality is in the form of

$$(2.16) \quad (R, E_f(\sqrt{\mu}) E_g(\sqrt{\mu})) = -\frac{1}{2\pi i} \int_{\gamma} m_1(\mu) E_f(\sqrt{\mu}) E_g(\sqrt{\mu}) d\mu,$$

by the formulas (2.8) and (2.15). From (2.12), (2.14) and Theorem.2.4, $m_1(\mu)$ can be extended as a finite meromorphic function from outside the strip (2.4) onto $\mathbf{C} \setminus \mathbf{R}_+$. By Theorem 2.4 we have that L has a finite number of spectral singularities; μ_1, \dots, μ_k , and eigenvalues; μ_{k+1}, \dots, μ_m each of them having finite multiplicity $\alpha_1, \dots, \alpha_k$ and $\alpha_{k+1}, \dots, \alpha_m$, respectively. From (2.14), if (1.4) holds $m_1(\mu)$ is a meromorphic function in $\mathbf{C} \setminus \mathbf{R}_+$ and has a finite number of poles μ_1, \dots, μ_m with finite multiplicities. In the neighborhood of the spectral singularities μ_1, \dots, μ_k , $m_1(\mu)$ satisfies the conditions

$$(2.17) \quad |m_1(u + iv)| \leq \frac{c_i}{|v|^{\alpha_i}}$$

for $a_i \leq u \leq a_{i+1}$, where $\mu = u + iv$, α_i , $0 < \alpha_i < \infty$, is the multiplicity of μ_i , $i = 1, \dots, k$, and

$$(2.18) \quad \bigcup_{i=1}^k [a_i, a_{i+1}] = [0, N].$$

3. Integral Representation of $m_1(\mu)$.

To derive an expansion formula in terms of eigenfunctions of the operator L , for (2.16), we will transform the integration over γ enclosing the spectrum, to the integration over spectrum. In order to do this we need the integral representation of the function $m_1(\mu)$ similar to that of given in [1].

Let us suppose that the following conditions hold:

- (a) the operator L has no eigenvalues,
- (b) there exists only one spectral singularity of the operator L with finite multiplicity α .

Clearly $m_1(\mu)$ is an analytic function in $\mathbf{C} \setminus \mathbf{R}_+$ and in the neighborhood of the spectral singularity, the following condition holds

$$(3.1) \quad |m_1(\mu)| \leq \frac{c}{|\operatorname{Im} \mu|^\alpha},$$

for c is a constant and $0 < \alpha < \infty$.

Lemma 3.1. *If (a) and (b) are satisfied, then there exist continuous complex functions $\sigma_1(t)$ and $\sigma_2(t)$ defined on the segment $[0, N]$ and a function $n(\mu)$ analytic in $\mathbf{C} \setminus \mathbf{R}_+$, continuous up to the boundary for which*

$$(3.2) \quad m_1(\mu) = \int_0^N \frac{\sigma_1(t)}{\mu - t} dt + (\alpha + 1)! \int_0^N \frac{\sigma_2(t)}{(\mu - t)^{\alpha+2}} dt + n(\mu).$$

Proof. Let us consider the following functions ;

$$(3.3) \quad M(\mu) = \begin{cases} m_1(\mu); & \mu \in [0, N] \\ 0; & \mu \in (N, \infty) \\ \frac{m_1(\mu)}{2}; & \mu \in \mathbf{C} \setminus \mathbf{R}_+ \end{cases},$$

$$(3.4) \quad n(\mu) = \begin{cases} 0; & \mu \in [0, N] \\ m_1(\mu); & \mu \in (N, \infty) \\ \frac{m_1(\mu)}{2}; & \mu \in \mathbf{C} \setminus \mathbf{R}_+ \end{cases}.$$

Therefore we can write the function $m_1(\mu)$ as follows:

$$(3.5) \quad m_1(\mu) = M(\mu) + n(\mu),$$

here $M(\mu)$ and $n(\mu)$ tend to zero as $|\mu|$ tends to infinity by (2.6), moreover $M(\mu)$ is analytic in $C \setminus [0, N]$ and in the neighborhood of the spectral singularity; $|M(\mu)| \leq \frac{c}{|\operatorname{Im} \mu|^\alpha}$ holds, here $c > 0$ is a constant, $0 < \alpha < \infty$. Therefore using the method similar to that of [1], we obtain the representation (3.2).

Now we assume that the following condition holds.

(c) the operator L has a finite number of spectral singularities each of them is of finite multiplicity.

Lemma 3.2. *If the conditions (a) and (c) are satisfied, then there exist continuous complex functions $\rho_i(t)$, defined on $[a_i, a_{i+1}]$, $i = 1, \dots, k$, respectively, a piecewise continuous function $\sigma(t)$ defined on $[0, N]$ and a function $n(\mu)$ analytic in $C \setminus \mathbf{R}_+$, continuous up to the boundary such that*

$$(3.6) \quad m_1(\mu) = \int_0^N \frac{\sigma(t)}{\mu - t} dt + \sum_{i=1}^k (\alpha_i + 1)! \int_{a_i}^{a_{i+1}} \frac{\rho_i(t)}{(\mu - t)^{\alpha_i + 2}} dt + n(\mu).$$

Proof. From (3.3) – (3.5) proof can be obtained easily.

Finally we assume that the following condition holds.

(d) the operator L has a finite number of spectral singularities $\mu_i \in [a_i, a_{i+1}]$, with multiplicities α_i , $0 < \alpha_i < \infty$, $i = 1, \dots, k$, and a finite number of eigenvalues μ_j , with multiplicities α_j , $0 < \alpha_j < \infty$, $j = k + 1, \dots, m$, respectively.

Now we arrive at the following.

Theorem 3.3. *If the condition (1.4) holds, then there exists continuous complex functions $\rho_i(t)$ defined on $[a_i, a_{i+1}]$, $i = 1, \dots, k$, a piecewise continuous complex function $\sigma(t)$ defined on $[0, N]$ and a function $n(\mu)$ analytic in $C \setminus \mathbf{R}_+$, continuous up to the boundary such that*

$$(3.7) \quad m_1(\mu) = \frac{1}{\Phi(\mu)} \int_0^N \frac{\sigma(t)}{\mu - t} dt + \sum_{i=1}^k \frac{(\alpha_i + 1)!}{\Phi(\mu)} \int_{a_i}^{a_{i+1}} \frac{\rho_i(t)}{(\mu - t)^{\alpha_i + 2}} dt + \frac{n(\mu)}{\Phi(\mu)},$$

where

$$(3.8) \quad \Phi(\mu) = \begin{cases} \prod_{j=k+1}^m (\mu - \mu_j)^{\alpha_j}; & \mu \in C_R \\ 1; & \mu \in \mathbf{C} \setminus C_R \end{cases},$$

and $C_R = \{\mu : |\mu| \leq R\}$ so that $\mu_j \in C_R$, $j = k + 1, \dots, m$.

Proof. It is clear that under the condition (1.4) the operator L satisfies (d). From Lemma (3.1) and Lemma (3.2) we reach to the desired result.

4. Expansion in Terms of the Principal Functions

We will call the functions $w(x, \sqrt{\mu})$, $\mu > 0$, the principal functions of the continuous spectrum. Note that these functions do not belong to $L^2(\mathbf{R}_+)$. In order to find an expansion formula for finite functions in terms of principal functions of the operator L , we introduce the principal functions of the spectral singularities by

$$(4.1) \quad w^{(s)}(x, \sqrt{\mu_i}) = \left(\frac{d}{d\mu}\right)^s \{w(x, \sqrt{\mu})\}_{\mu=\mu_i},$$

$s = 0, 1, \dots, \alpha_i - 1$, $i = 1, \dots, k$, and principal functions of the discrete spectrum by

$$(4.2) \quad w^{(s)}(x, \sqrt{\mu_i}) = \left(\frac{d}{d\mu}\right)^s \{w(x, \sqrt{\mu})\}_{\mu=\mu_i} \in L^2(\mathbf{R}_+),$$

$s = 0, 1, \dots, \alpha_i - 1$, $i = k + 1, \dots, m$, where $w(x, \sqrt{\mu_i})$ is an eigenfunction and $w^{(1)}(x, \sqrt{\mu_i}), \dots, w^{(\alpha_i-1)}(x, \sqrt{\mu_i})$ are associated functions corresponding to the eigenvalues μ_i , $i = k + 1, \dots, m$. Obviously $w^{(s)}(x, \sqrt{\mu_i}) \notin L^2(\mathbf{R}_+)$, $i = 1, \dots, k$, $s = 0, 1, \dots, \alpha_i - 1$.

Remark 4.1. From Theorem (2.3) we easily get that

$$\sup_{x \in \mathbf{R}_+} \frac{|w^{(s)}(x, \sqrt{\mu})|}{(1+x)^s} < \infty, \quad \mu > 0, \quad s = 0, 1, 2, \dots,$$

here $w^{(s)} = \left(\frac{\partial}{\partial \mu}\right)^{(s)} w$.

Let us recall the Hilbert spaces ([6])

$$H_n = \left\{ f : \int_0^\infty (1+x)^{2n} |f(x)|^2 dx < \infty \right\},$$

$$H_{-n} = \left\{ g : \int_0^\infty (1+x)^{-2n} |g(x)|^2 dx < \infty \right\}$$

with respective norms

$$\|f\|_n^2 = \int_0^\infty (1+x)^{2n} |f(x)|^2 dx, \quad \|g\|_{-n}^2 = \int_0^\infty (1+x)^{-2n} |g(x)|^2 dx,$$

$n = 0, 1, \dots$. Clearly

$$H_0 = L^2(\mathbf{R}_+), \quad H_n L^2(\mathbf{R}_+) H_{-n}, \quad n = 1, 2, \dots$$

It is known that H_{-n} is isomorphic to the dual of H_n ($H_{-n} \sim H'_n$).

Let us choose s_0 such that $s_0 = \max\{\alpha_1, \dots, \alpha_k\}$. In this case

$$(4.3) \quad H_{(s_0+2)} \subsetneq L^2(\mathbf{R}_+) \subsetneq H_{-(s_0+2)}.$$

We obtain for the principal functions of the spectral singularities that

$$\frac{\partial^s}{\partial \mu^2} \{w(x, \sqrt{\mu})\}_{\mu=\mu_i} \in H_{-(s+1)}$$

for $s = 0, 1, \dots, \alpha_i - 1$, $i = 1, \dots, k$, hence, we get that

$$\frac{\partial^s}{\partial \mu^2} \{w(x, \sqrt{\mu})\}_{\mu=\mu_i} \in H_{-s_0}$$

for $s = 0, 1, \dots, \alpha_i - 1$, $i = 1, \dots, k$, by the choice of s_0 .

Now, we need the following two lemmas ([6]).

Lemma 4.2. *There is a number $c > 0$ such that for every finite function $f \in L^2(\mathbf{R}_+)$*

$$\int_0^\infty |E_f(\sqrt{\mu})|^2 \sqrt{\mu} d\mu \leq c \int_0^\infty |f(x)|^2 dx$$

holds.

Lemma 4.3. *If*

$$\int_0^\infty (1+x)^{2n} |f(x)|^2 dx < \infty$$

holds, then the function $E_f(\sqrt{\mu})$ has a derivative of order $(n-1)$ which is absolutely continuous on every finite interval of the half axis $\mu > 0$. There exists a number $c_n > 0$ such that

$$\int_{-\infty}^{\infty} \left| \left(\frac{d}{d\lambda} \right)^n \lambda E_f(\lambda) \right|^2 d\lambda \leq c_n \int_0^{\infty} (1+x)^{2n} |f(x)|^2 dx$$

holds.

Let us suppose the following notations:

$$H_{\pm} := H_{\pm(s_0+2)},$$

$$(4.4) \quad \Phi_j(\mu) := \prod_{l=k+1, j \neq l}^m (\mu - \mu_l)^{\alpha_l},$$

$j = k+1, \dots, m, \mu \in \mathbf{C} \setminus \mathbf{R}_+$.

$$(4.5) \quad \tilde{w}^{(\alpha_i+1-s_i)}(x, \sqrt{t}) := \left(\frac{d}{d\mu} \right)^{(\alpha_i+1-s_i)} \left\{ \frac{w(x, \sqrt{\mu})}{\Phi(\mu)} \right\}_{\mu=t},$$

$$(4.6) \quad \psi^{(\alpha_j-1-s_j)}(t, \mu_j; x) := \left(\frac{d}{d\mu} \right)^{\alpha_j-1-s_j} \left\{ \frac{w(x, \sqrt{\mu})}{(\mu-t)\Phi_j(\mu)} \right\}_{\mu=\mu_j},$$

$$\tilde{\psi}^{(\alpha_j-1-s_j)}(t, \mu_j, \alpha_i; x)$$

$$(4.7) \quad := \left(\frac{d}{d\mu} \right)^{(\alpha_j-1-s_j)} \left\{ \frac{w(x, \sqrt{\mu})}{(\mu-t)^{\alpha_i+2} \Phi_j(\mu)} \right\}_{\mu=\mu_j},$$

$i = 1, \dots, k, j = k+1, \dots, m$, where Φ is the function defined by (3.8).

$$(4.8) \quad (f, \tilde{w}^{(\alpha_i+1-s_i)}(\sqrt{t})) := \int_0^{\infty} f(y) \tilde{w}^{(\alpha_i+1-s_i)}(y, \sqrt{t}) dy,$$

$$(f, \psi^{(\alpha_j-1-s_j)}(t, \mu_j)) := \int_0^{\infty} f(y) \psi^{(\alpha_j-1-s_j)}(t, \mu_j; y) dy,$$

$$(f, \tilde{\psi}^{(\alpha_j-1-s_j)}(t, \mu_j; \alpha_i)) := \int_0^{\infty} f(y) \tilde{\psi}^{(\alpha_j-1-s_j)}(t, \mu_j, \alpha_i; y) dy.$$

Theorem 4.4. Under the condition (1.4), the following spectral expansion formula in terms of the principal functions of L holds

$$\begin{aligned}
f(x) &= \int_0^N \frac{\sigma(t)}{\Phi(t)} w(x, \sqrt{t}) E_f(\sqrt{t}) dt \\
&+ \sum_{i=1}^k \sum_{s_i=0}^{\alpha_i+1} C_{\alpha_i+1}^{s_i} \int_{\alpha_i}^{\alpha_i+1} \rho_i(t) \left(f, \tilde{w}^{(\alpha_i+1-s_i)}(\sqrt{t}) \right) w^{(s_i)}(x, \sqrt{t}) dt \\
&+ \sum_{j=k+1}^m \sum_{s_j=0}^{\alpha_j-1} \frac{1}{(\alpha_j-1)!} C_{\alpha_j-1}^{s_j} \times \\
&\times \int_0^N \sigma(t) \left(f, \psi^{(\alpha_j-1-s_j)}(t, \mu_j) \right) w^{(s_j)}(x, \sqrt{t}) dt \\
(4.9) \quad &+ \sum_{i=1}^k \sum_{j=k+1}^m \sum_{s_j=0}^{\alpha_j-1} \frac{(\alpha_i+1)!}{(\alpha_j-1)!} C_{\alpha_j-1}^{s_j} \times \\
&\times \int_{\alpha_i}^{\alpha_i+1} \rho_i(t) \left(f, \tilde{\psi}^{(\alpha_j-1-s_j)}(t, \mu_j; \alpha_j) \right) w^{(s_j)}(x, \sqrt{t}) dt \\
&+ \frac{1}{2\pi i} \int_N^\infty \frac{1}{\Phi(\mu)} \{m_1(\mu+i0) - m_1(\mu-i0)\} w(x, \sqrt{\mu}) E_f(\sqrt{\mu}) d\mu
\end{aligned}$$

for all $f \in H_+$, where C_α^s are the binomial coefficients. The integrals in (4.9) converge in the norm of H_- .

Proof. Substituting the representation for $m_1(\mu)$ given by (3.7) in the Parseval equality (2.16), we obtain that

$$\begin{aligned}
(f, g) &= \frac{1}{2\pi i} \int_0^N \sigma(t) \int_\gamma \frac{E_f(\sqrt{\mu}) E_g(\sqrt{\mu})}{(\mu-t)\Phi(\mu)} d\mu dt \\
(4.10) \quad &+ \frac{1}{2\pi i} \sum_{i=1}^k (\alpha_i+1)! \int_{\alpha_i}^{\alpha_i+1} \rho_i(t) \int_\gamma \frac{E_f(\sqrt{\mu}) E_g(\sqrt{\mu})}{(\mu-t)^{\alpha_i+2} \Phi(\mu)} d\mu dt
\end{aligned}$$

$$+ \frac{1}{2\pi i} \int_N^{\infty} \{m_1(\mu + i0) - m_1(\mu - i0)\} \frac{E_f(\sqrt{\mu}) E_g(\sqrt{\mu}) d\mu}{\Phi(\mu)},$$

where

$$\begin{aligned} & \int_{\gamma}^{\infty} n(\mu) E_f(\sqrt{\mu}) E_g(\sqrt{\mu}) d\mu \\ &= \int_N^{\infty} \{m_1(\mu + i0) - m_1(\mu - i0)\} E_f(\sqrt{\mu}) E_g(\sqrt{\mu}) d\mu \end{aligned}$$

by (3.4). If we evaluate the contour integrals in (4.10), using Cauchy Residue theory and taking care of the notations (4.4) – (4.8), we get the expansion (4.9) for $f \in C_0^\infty \subset H_+$. Since C_0^∞ is dense in H_+ we have the expansion for $f \in H_+$.

Now we define the following operators;

$$T_0 f(x) \quad : \quad = \int_0^N \frac{\sigma(t)}{\Phi(t)} w(x, \sqrt{t}) E_f(\sqrt{t}) dt,$$

$$\begin{aligned} T_{\tilde{w}} f(x) &:= \sum_{i=1}^k \sum_{s_i=0}^{\alpha_i+1} C_{\alpha_i+1}^{s_i} \times \\ &\times \int_{a_i}^{a_i+1} \rho_i(t) \left(f, \tilde{w}^{(\alpha_i+1-s_i)}(\sqrt{t}) \right) w^{(s_i)}(x, \sqrt{t}) dt, \end{aligned}$$

$$\begin{aligned} T_{\psi} f(x) &:= \sum_{j=k+1}^m \sum_{s_j=0}^{\alpha_j-1} \frac{1}{(\alpha_j-1)!} C_{\alpha_j-1}^{s_j} \times \\ &\times \int_0^N \sigma(t) \left(f, \psi^{(\alpha_j-1-s_j)}(t, \mu_j) \right) w^{(s_j)}(x, \sqrt{t}) dt, \end{aligned}$$

$$\begin{aligned} T_{\tilde{\psi}} f(x) &:= \sum_{i=1}^k \sum_{j=k+1}^m \sum_{s_j=0}^{\alpha_j-1} \frac{(\alpha_j+1)!}{(\alpha_j-1)!} C_{\alpha_j-1}^{s_j} \times \\ &\times \int_{a_i}^{a_i+1} \rho_i(t) \left(f, \tilde{\psi}^{(\alpha_j-1-s_j)}(t, \mu_j; \alpha_j) \right) w^{(s_j)}(x, \sqrt{t}) dt, \end{aligned}$$

$$T_N f(x) := \frac{1}{2\pi i} \int_N^{\infty} \frac{1}{\Phi(\mu)} \{m_1(\mu + i0) - m_1(\mu - i0)\} w(x, \sqrt{\mu}) E_f(\sqrt{\mu}) d\mu.$$

We now show that each of the operators T_0 , $T_{\tilde{w}}$, T_ψ , $T_{\tilde{\psi}}$ and T_N is continuous from H_+ onto H_- .

T_0 is continuous from $L^2(\mathbf{R}_+)$ onto itself by Lemma 4.2. Using isomorphism

$H_{-s_0} \sim H'_{s_0}$ we get that T_0 is continuous from H_{s_0} onto H_{-s_0} or from H_+ onto H_- by inclusions (4.3). Hence there exists a constant c_0 such that

$$(4.11) \quad \|T_0 f\|_- \leq c_0 \|f\|_+ \quad 4.11$$

for any $f \in H_+$.

$T_{\tilde{w}}$ includes the principal functions of the continuous spectrum. Since

$$\begin{aligned} \|T_{\tilde{w}} f(x)\|_- &= \int_0^\infty \left| \frac{T_{\tilde{w}} f(x)}{(1+x)^{s_0+2}} \right|^2 dx \\ &\leq \sum_{i=1}^k \sum_{s_i=0}^{\alpha_i+1} C_{\alpha_i+1}^{s_i} \int_0^\infty \int_{a_i}^{a_{i+1}} \left| \frac{\rho_i(t) w^{(s_i)}(x, \sqrt{t})}{(1+x)^{s_0+2}} \right|^2 dt dx \times \\ &\quad \times \int_{a_i}^{a_{i+1}} \left| (f, \tilde{w}^{(\alpha_i+1-s_i)}(\sqrt{t})) \right|^2 dt. \end{aligned}$$

We see that first double integral is convergent from (4.5), Remark 4.1, Lemma 4.2 and Lemma 4.3. Therefore we get from Lemma 4.2 and Lemma 4.3 that

$$(4.12) \quad \|T_{\tilde{w}} f\|_- \leq c_{\tilde{w}} \|f\|_+,$$

where $c_{\tilde{w}} > 0$ is a constant.

T_ψ and $T_{\tilde{\psi}}$ are continuous from $L^2(\mathbf{R}_+)$ onto itself by (4.2), (4.6) and (4.7), so we can write that

$$\|T_\psi\| \leq c_\psi \|f\|, \quad \|T_{\tilde{\psi}}\| \leq c_{\tilde{\psi}} \|f\|$$

where $c_\psi > 0$, and $c_{\tilde{\psi}} > 0$ are constants. Using (4.3) and isomorphism $H_{-s_0} \sim H'_{s_0}$ we get that T_ψ and $T_{\tilde{\psi}}$ are continuous from H_+ onto H_- . Hence we have

$$(4.13) \quad \|T_\psi f\|_- \leq c_\psi \|f\|_+, \quad \|T_{\tilde{\psi}} f\|_- \leq c_{\tilde{\psi}} \|f\|_+$$

for the constants given above.

Finally we consider the operator T_N . From (2.10) and Theorem 2.3, using a similar proof to that of the well-known Lemma 4.2 we get that

$$\|T_N f\| \leq c_N \|f\|,$$

here $c_N > 0$ is a constant. From the inclusions (4.3) and the isomorphism we find that

$$(4.14) \quad \|T_N f\|_- \leq c_N \|f\|_+,$$

where $c_N > 0$ is a constant.

We get the proof of the theorem by (4.11), (4.12), (4.13) and (4.14).

References

- [1] E. Bairamov and A. O. Çelebi, Spectrum and spectral expansion for the non-selfadjoint discrete Dirac operators, *Quart. J. Math. Oxford Ser.*, (2), **50**, pp. 371–384, (1999).
- [2] E. Bairamov, Ö. Çakar and A. O. Çelebi, Quadratic pencil of Schrödinger operators with spectral singularities: discrete spectrum and principal functions, *J. Math. Anal. Appl.*, **216**, No. 1, pp. 303–320, (1997).
- [3] Lyance, A differential Operator with Spectral Singularities, I,II, *Amer. Math. Soc. Trans. Ser. 2*, Vol. **60**, 185–225, pp. 227–283, (1967).
- [4] V. A. Marchenko, Expansion in eigenfunctions of Non-selfadjoint Singular Second-Order Differential Operators, *A MS Translations*, (2) **25**, pp. 77–130, (1963).
- [5] V. A. Marchenko, Sturm-Liouville operators and applications Operator Theory: Advances and Applications, 22. Birkhäuser Verlag, Basel, (1986).
- [6] M.A.Naimark, Linear Differential Operators I-II, Ungar, New-York, (1968).
- [7] J.T.Schwartz, Some non-selfadjoint operators, *Comm. Pure and Appl. Math.* **13**, pp. 609–639, (1960).

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